

4.7 High energy and nuclear physics

Nuclear decay

Nuclear decay law	$N(t) = N(0)e^{-\lambda t}$	(4.163)	$N(t)$ number of nuclei remaining after time t t time
Half-life and mean life	$T_{1/2} = \frac{\ln 2}{\lambda}$	(4.164)	λ decay constant
	$\langle T \rangle = 1/\lambda$	(4.165)	$T_{1/2}$ half-life $\langle T \rangle$ mean lifetime
Successive decays $1 \rightarrow 2 \rightarrow 3$ (species 3 stable)	$N_1(t) = N_1(0)e^{-\lambda_1 t}$	(4.166)	N_1 population of species 1
	$N_2(t) = N_2(0)e^{-\lambda_2 t} + \frac{N_1(0)\lambda_1(e^{-\lambda_1 t} - e^{-\lambda_2 t})}{\lambda_2 - \lambda_1}$	(4.167)	N_2 population of species 2
	$N_3(t) = N_3(0) + N_2(0)(1 - e^{-\lambda_2 t}) + N_1(0)\left(1 + \frac{\lambda_1 e^{-\lambda_2 t} - \lambda_2 e^{-\lambda_1 t}}{\lambda_2 - \lambda_1}\right)$	(4.168)	N_3 population of species 3 λ_1 decay constant $1 \rightarrow 2$ λ_2 decay constant $2 \rightarrow 3$
Geiger's law ^a	$v^3 = a(R - x)$	(4.169)	v velocity of α particle x distance from source a constant
Geiger–Nuttall rule	$\log \lambda = b + c \log R$	(4.170)	R range b, c constants for each series α , β , and γ

^aFor α particles in air (empirical).

Nuclear binding energy

Liquid drop model ^a	$B = a_v A - a_s A^{2/3} - a_c \frac{Z^2}{A^{1/3}} - a_a \frac{(N - Z)^2}{A} + \delta(A)$	(4.171)	N number of neutrons A mass number ($= N + Z$) B semi-empirical binding energy Z number of protons a_v volume term ($\sim 15.8 \text{ MeV}$) a_s surface term ($\sim 18.0 \text{ MeV}$) a_c Coulomb term ($\sim 0.72 \text{ MeV}$) a_a asymmetry term ($\sim 23.5 \text{ MeV}$) a_p pairing term ($\sim 33.5 \text{ MeV}$)
	$\delta(A) \simeq \begin{cases} +a_p A^{-3/4} & Z, N \text{ both even} \\ -a_p A^{-3/4} & Z, N \text{ both odd} \\ 0 & \text{otherwise} \end{cases}$	(4.172)	
Semi-empirical mass formula	$M(Z, A) = Z M_H + N m_n - B$	(4.173)	$M(Z, A)$ atomic mass M_H mass of hydrogen atom m_n neutron mass

^aCoefficient values are empirical and approximate.



Nuclear collisions

Breit–Wigner formula ^a	$\sigma(E) = \frac{\pi}{k^2} g \frac{\Gamma_{ab}\Gamma_c}{(E - E_0)^2 + \Gamma^2/4} \quad (4.174)$	$\sigma(E)$ cross-section for $a + b \rightarrow c$
	$g = \frac{2J + 1}{(2s_a + 1)(2s_b + 1)} \quad (4.175)$	k incoming wavenumber g spin factor E total energy (PE + KE) E_0 resonant energy Γ width of resonant state R
Total width	$\Gamma = \Gamma_{ab} + \Gamma_c \quad (4.176)$	Γ_{ab} partial width into $a + b$ Γ_c partial width into c
Resonance lifetime	$\tau = \frac{\hbar}{\Gamma} \quad (4.177)$	τ resonance lifetime J total angular momentum quantum number of R
Born scattering formula ^b	$\frac{d\sigma}{d\Omega} = \left \frac{2\mu}{\hbar^2} \int_0^\infty \frac{\sin Kr}{Kr} V(r) r^2 dr \right ^2 \quad (4.178)$	$s_{a,b}$ spins of a and b $\frac{d\sigma}{d\Omega}$ differential collision cross-section μ reduced mass $K = \mathbf{k}_{in} - \mathbf{k}_{out} $ (see footnote) r radial distance $V(r)$ potential energy of interaction
Mott scattering formula ^c	$\frac{d\sigma}{d\Omega} = \left(\frac{\alpha}{4E} \right)^2 \left[\csc^4 \frac{\chi}{2} + \sec^4 \frac{\chi}{2} + \frac{A \cos \left(\frac{\alpha}{\hbar v} \ln \tan^2 \frac{\chi}{2} \right)}{\sin^2 \frac{\chi}{2} \cos \frac{\chi}{2}} \right] \quad (4.179)$	$\hbar$ (Planck constant)/ 2π α/r scattering potential energy χ scattering angle v closing velocity $A = 2$ for spin-zero particles, $= -1$ for spin-half particles
	$\frac{d\sigma}{d\Omega} \simeq \left(\frac{\alpha}{2E} \right)^2 \frac{4 - 3 \sin^2 \chi}{\sin^4 \chi} \quad (A = -1, \alpha \ll v\hbar) \quad (4.180)$	

^aFor the reaction $a + b \leftrightarrow R \rightarrow c$ in the centre of mass frame.

^bFor a central field. The Born approximation holds when the potential energy of scattering, V , is much less than the total kinetic energy. K is the magnitude of the change in the particle's wavevector due to scattering.

^cFor identical particles undergoing Coulomb scattering in the centre of mass frame. Nonidentical particles obey the Rutherford scattering formula (page 72).

Relativistic wave equations^a

Klein–Gordon equation (massive, spin zero particles)	$(\nabla^2 - m^2)\psi = \frac{\partial^2 \psi}{\partial t^2} \quad (4.181)$	ψ wavefunction m particle mass t time
Weyl equations (massless, spin 1/2 particles)	$\frac{\partial \psi}{\partial t} = \pm \left(\sigma_x \frac{\partial \psi}{\partial x} + \sigma_y \frac{\partial \psi}{\partial y} + \sigma_z \frac{\partial \psi}{\partial z} \right) \quad (4.182)$	ψ spinor wavefunction σ_i Pauli spin matrices (see page 26)
Dirac equation (massive, spin 1/2 particles)	$(\mathbf{i}\gamma^\mu \partial_\mu - m)\psi = 0 \quad (4.183)$	$\mathbf{i}$ $\mathbf{i}^2 = -1$ γ^μ Dirac matrices:
	where $\partial_\mu = \left(\frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \quad (4.184)$	$\gamma^0 = \begin{pmatrix} \mathbf{1}_2 & 0 \\ 0 & -\mathbf{1}_2 \end{pmatrix}$
	$(\gamma^0)^2 = \mathbf{1}_4; \quad (\gamma^1)^2 = (\gamma^2)^2 = (\gamma^3)^2 = -\mathbf{1}_4 \quad (4.185)$	$\gamma^i = \begin{pmatrix} 0 & \sigma_i \\ -\sigma_i & 0 \end{pmatrix}$ $\mathbf{1}_n$ $n \times n$ unit matrix

^aWritten in natural units, with $c = \hbar = 1$.

